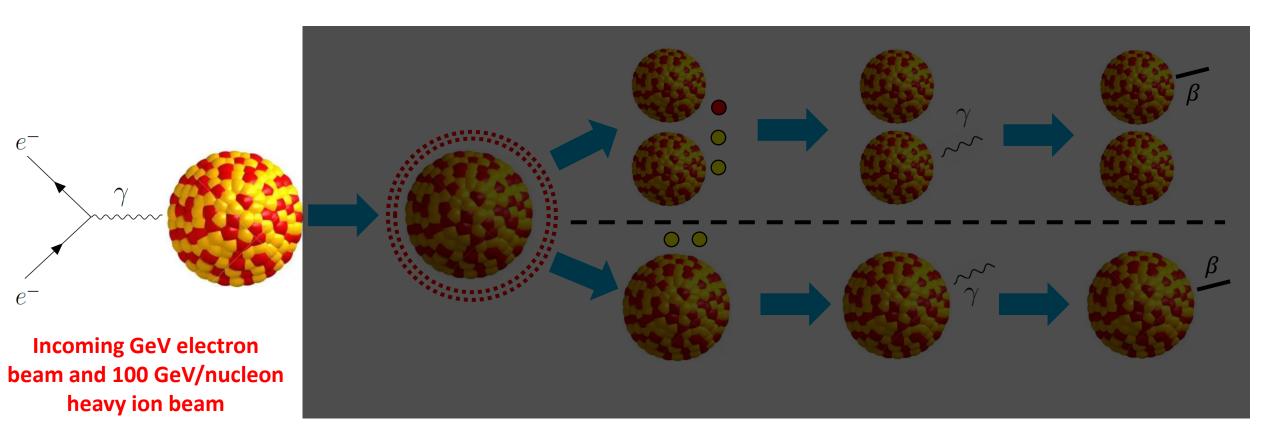
Production and detection of **exotic nuclei** at the EIC

Barak Schmookler,

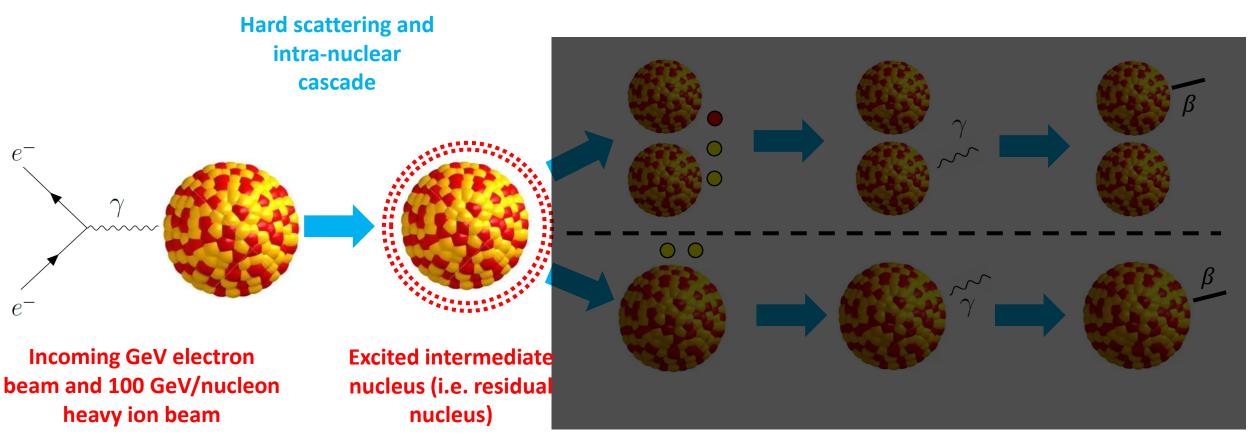
Ben Collis, Abhay Deshpande, Zach Finger,
Ciprian Gal, Mark Harvey, Brynna Moran, Pawel Nadel-Turonski

Motivating questions

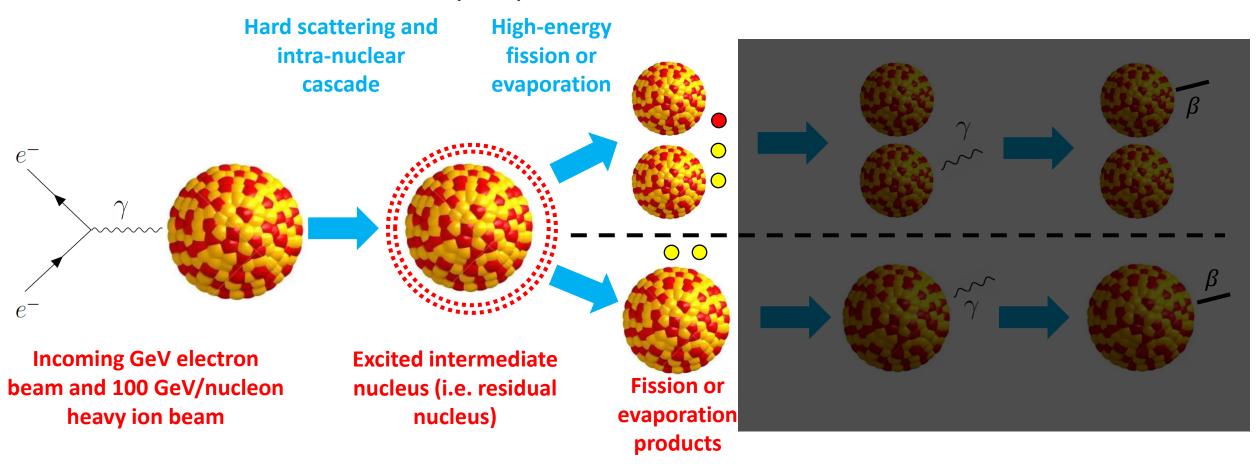
- □Can we use high-energy electron-heavy nucleus scattering at the future EIC to produce exotic nuclei (i.e. undiscovered rare isotopes)?
- □Can we go on to detect and correctly identify the produced exotic nuclei? Can we also study the level structure of the nuclei by detecting the decay photons? What requirements does this place on the far-forward detection area?
- □ If we can produce, detect, and identify exotic nuclei at the EIC, how can these results complement the work being done at dedicated rare isotope facilities?



t = 0



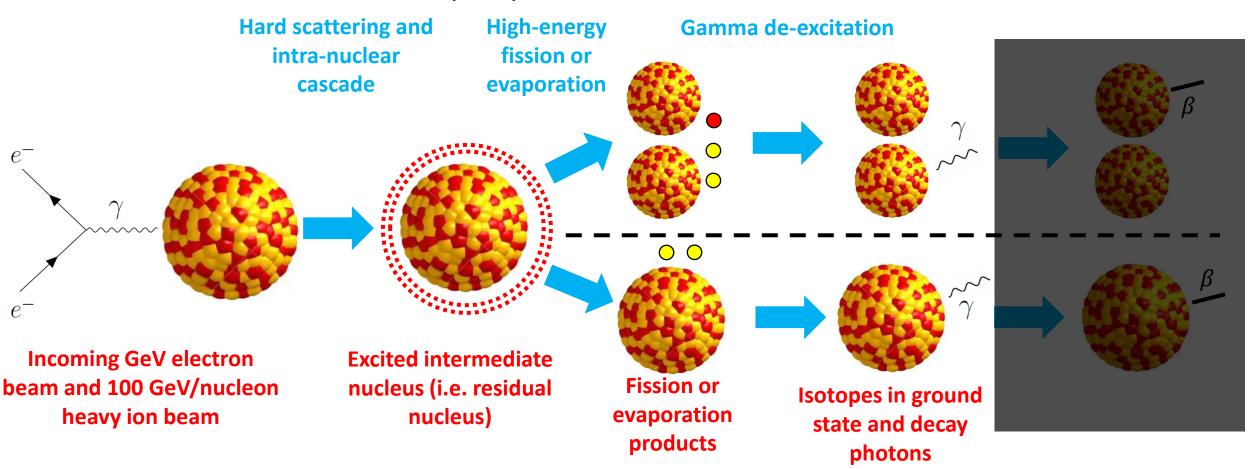
$$t = 0$$
 $t = 10^{-22} s$



$$t = 0$$

$$t = 10^{-22} s$$

$$t = 10^{-20} - 10^{-17} s$$

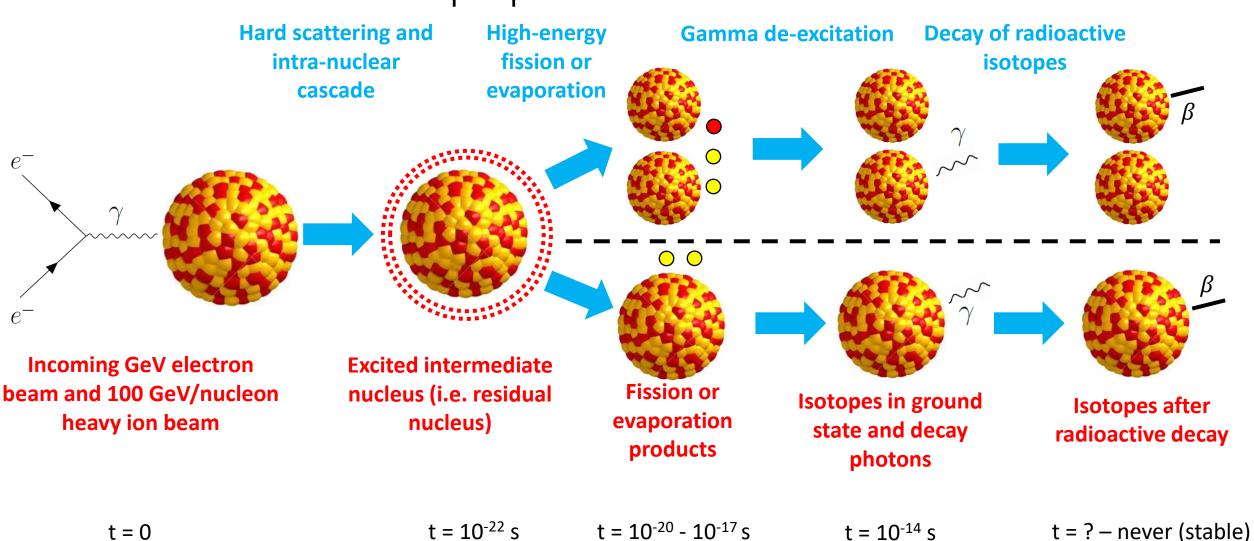


t = 0

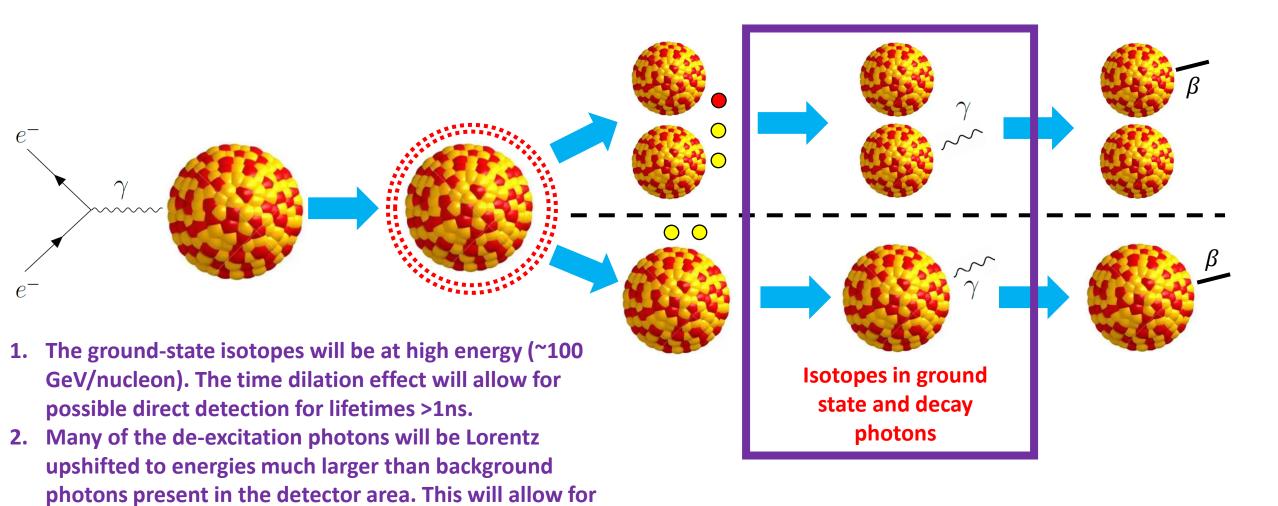
 $t = 10^{-22} s$

 $t = 10^{-20} - 10^{-17} s$

 $t = 10^{-14} s$



Where the EIC can potentially contribute

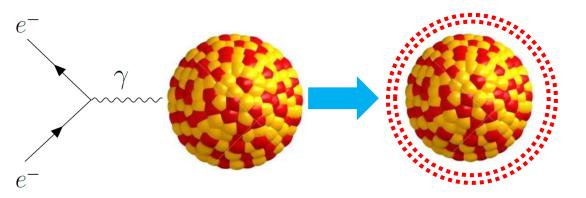


clean detection/identification of these photons, which

can be used to study the level-structure of the isotopes.

How can we study this?

Hard scattering and intra-nuclear cascade



Incoming GeV electron beam and 100 GeV/nucleon heavy ion beam

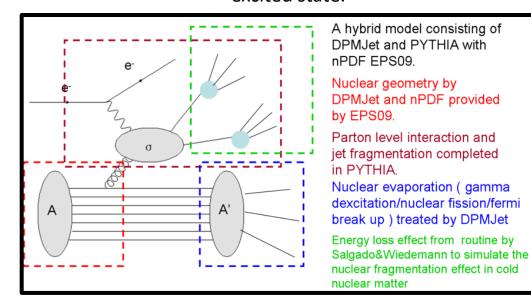
Excited intermediate nucleus (i.e. residual nucleus)

<u>Step 1</u>

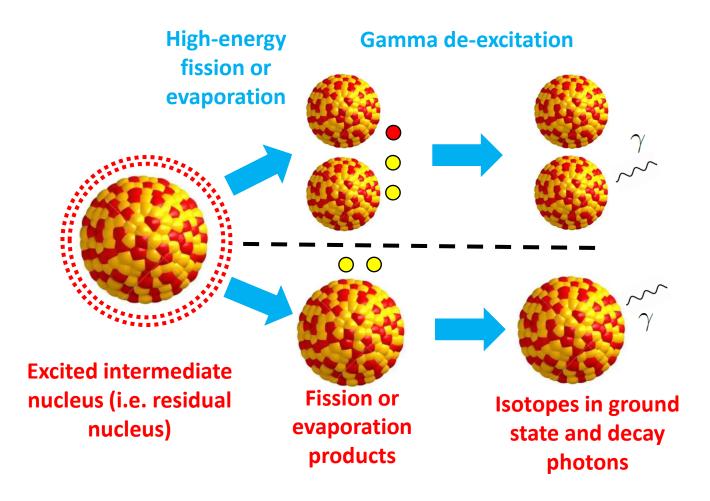
The hard scattering and the intra-nuclear cascade are modelled using the *Benchmark eA Generator* for Leptoproduction – BeAGLE

(https://wiki.bnl.gov/eic/index.php/BeAGLE)

This leaves us with the residual nucleus in an excited state.



How can we study this?



Step 2

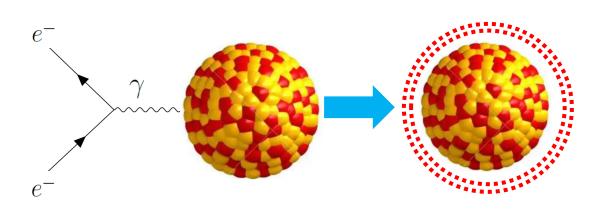
For each event, the residual nucleus with a given A, Z, and excitation energy is then handed over to either *FLUKA* (Annals of Nuclear Energy 82, 10-18 (2015)) or <u>ABLA07</u> for decay (fission or evaporation) followed by gamma de-excitation. We are left with the decay products of the residual nucleus.

FLUKA is used extensively in high-energy physics but has not been used for the study of rare isotope production. ABLAO7 is used extensively in the rare isotope community, and is the second part of the abrasion-ablation code ABRABLAO7. We run the BeAGLE events though both these codes and compare the results.

(N.B. The *FLUKA* decay of the residual nucleus has been directly incorporated into the *BeAGLE* simulation framework, allowing for easier analysis.

- Using *BeAGLE*, we simulate an 18 GeV electron beam colliding with a 110 GeV/nucleon ²³⁸U or ²⁰⁸Pb beam.
- ☐ We then study the excited residual nucleus that get created following the hard scattering and intra-nuclear cascade.
- The only relevant quantities are the A and Z of the residual, as well as its excitation energy.

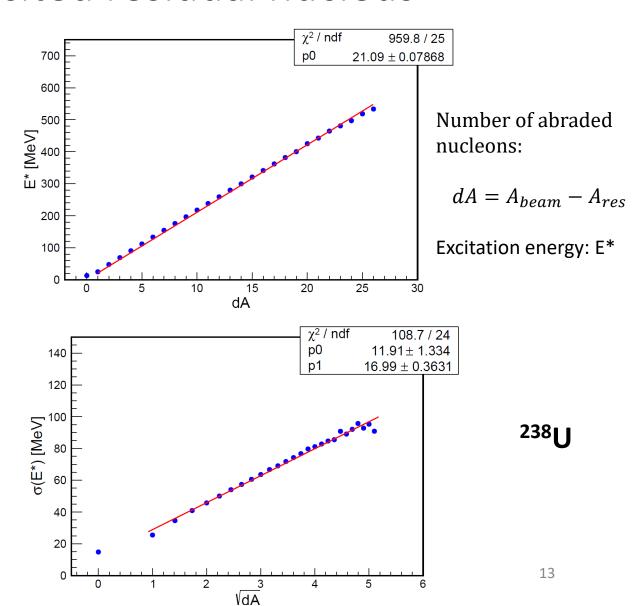
Hard scattering and intra-nuclear cascade



Excited intermediate nucleus (i.e. residual nucleus)

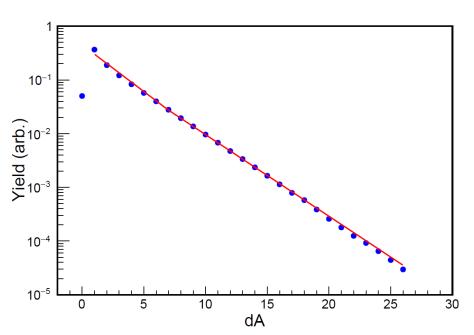
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 - The excitation energy shows a simple dependence on the number of abraded nucleons.



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- The cross section for abrading a given number of nucleons (for dA>1) shows a (piecewise) exponential dependence.



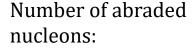
Number of abraded nucleons:

$$dA = A_{beam} - A_{res}$$

Excitation energy: E*

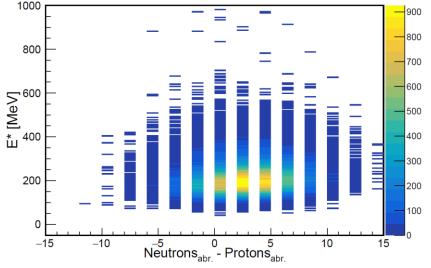
²³⁸U

- ■We find that the production of the residual nucleus in *BeAGLE* manifests as a very simple abrasion model:
 - The excitation energy shows a simple dependence on the number of abraded nucleons.
 - The cross section for abrading a given number of nucleons (for dA>1) shows a (piecewise) exponential dependence.
 - For a given number of abraded nucleons, the relative proportion of neutrons and protons abraded is close to a hypergeometric distribution



$$dA = A_{beam} - A_{res}$$

Excitation energy: E*

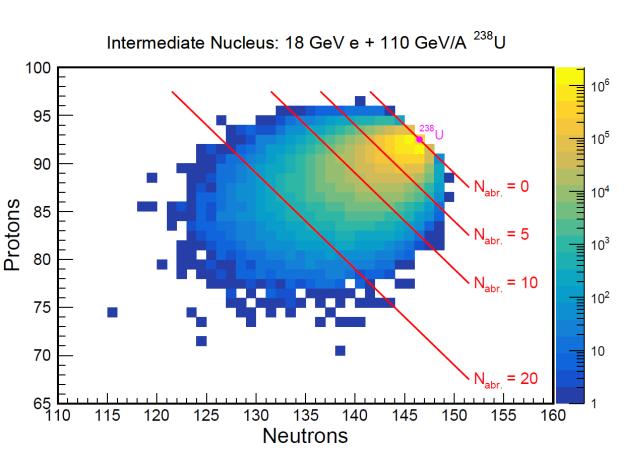


Excitation Energy vs. abrasion asymmetry: 10 nucleons abraded

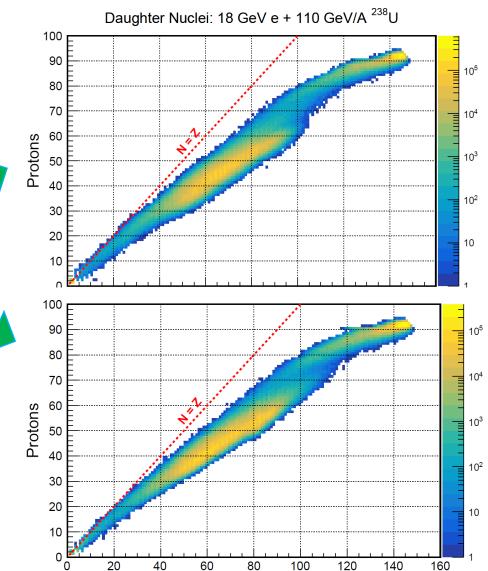
238_U

We can then decay the excited residual nucleus

ABLAO?

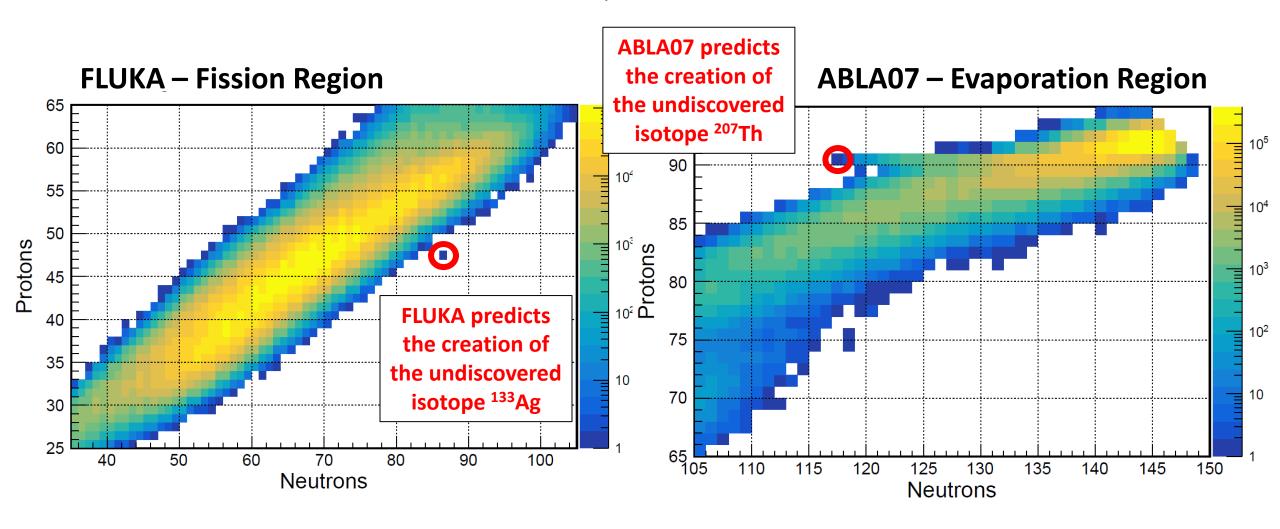


10 million events simulated

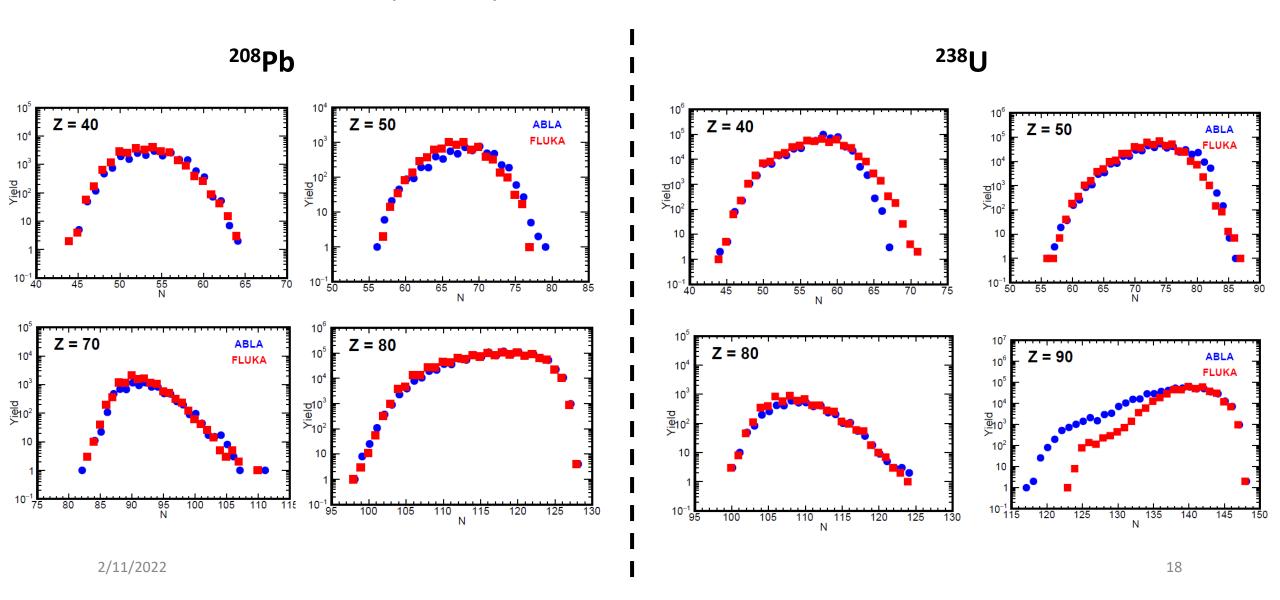


Neutrons

Using this 10 million ²³⁸U event sample, we see hints of exotic nuclei production



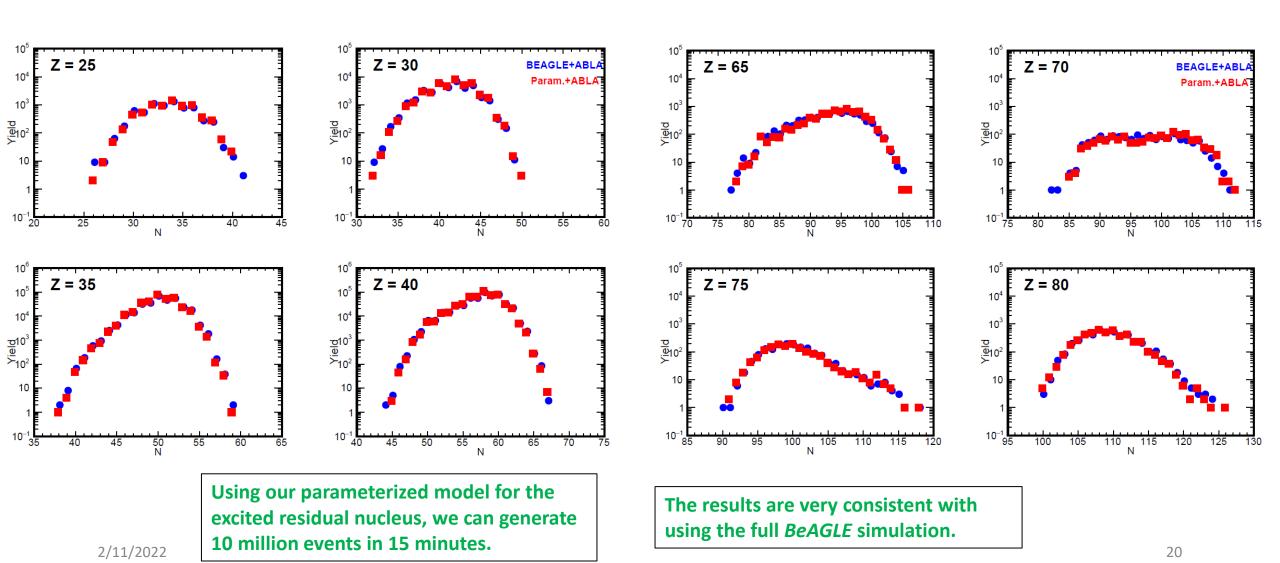
We can directly compare the results of FLUKA and ABLA07



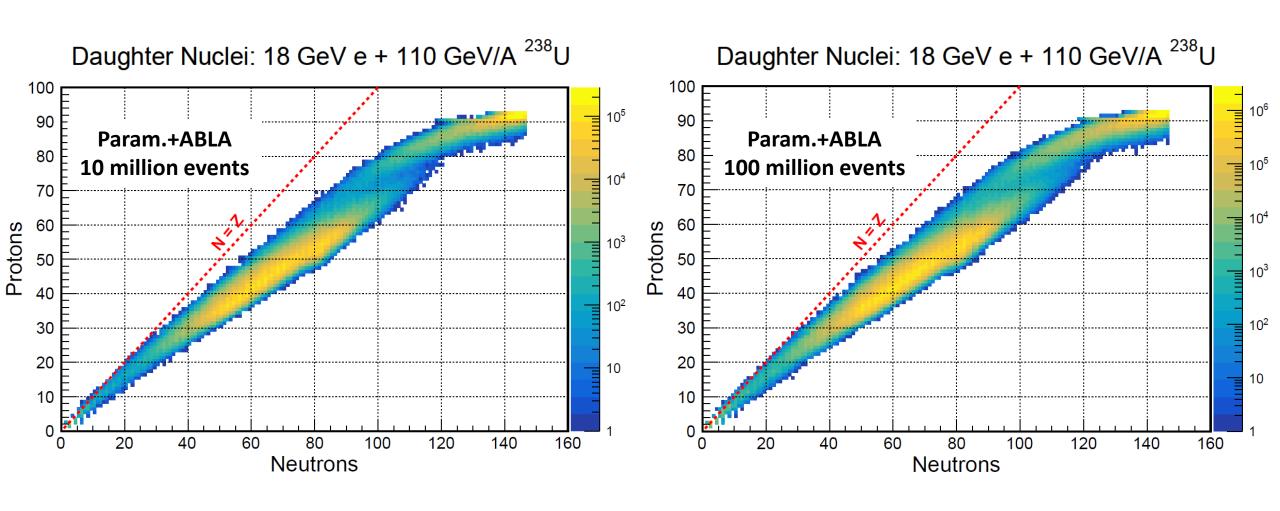
We need to simulate much more than 10 million events

☐ Using a *pythia6* simulation where we generate over the entire allowed kinematic phase space (i.e. all the way down to photoproduction), we see that the total scattering cross section is about 100 µb. \square If we make the assumptions that 1) we collect 10 fb⁻¹ integrated luminosity per year and 2) the production of nuclear isotopes is independent of the kinematics (i.e. Q^2 and x), we can estimate than 10 million events will correspond to about 5 minutes of running. □ Even though the above calculation is very rough, it clearly shows that we need to generate much larger quantities of events in order to study the EIC's capacity to produce rare isotopes. ☐Generating larger number of events with *BeAGLE* becomes computationally expensive. Fortunately for us, all we care about is the production of the excited residual nucleus. And we can create a simple parameterization of the residual nucleus production based on the *BeAGLE* model.

Comparison of *BeAGLE* results and parameterized distribution

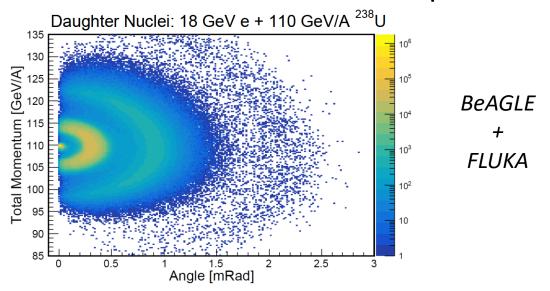


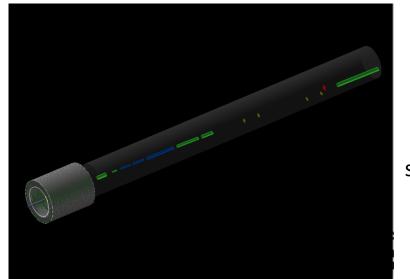
Towards higher statistics simulations



Detection and identification of the nuclear isotopes

- Dour simulation studies suggest that the daughter isotopes of the residual nucleus will be produced with (per-nucleon) momenta close the incoming ion beam momentum and with a very small scattering angles with respect to the ion beam.
- ☐ In order to detect these produced isotopes, we will need to use the *far forward* part of the interaction region.





Far forward magnets and detectors of IR6 in the *Fun4All* simulation framework

Isotope detection under a simple assumption

In the simplest assumption, the momentum pernucleon of the outgoing isotope (p_N) is the same as the momentum per-nucleon of the incoming beam $(p_{N,beam})$.

$$x_{L} = \frac{R}{R_{beam}} = \left[\frac{Ap_{N}}{Z} \right] / \frac{A_{beam}p_{N,beam}}{Z_{beam}}$$

$$= \left[\frac{A}{Z} \right] / \frac{A_{beam}p_{N,beam}}{Z_{beam}}$$

We can calculate various quantities — such as the isotope hit position at a Roman Pot (RP) and whether the isotope is within the RP acceptance — using the above equation.

Some definitions:

$$Rigidity = R = \frac{p}{Z}$$

$$x_L = \frac{R}{R_{heam}}$$

Relative Rigidity =
$$R_{Rel} = \frac{R - R_{beam}}{R_{beam}} = x_L - 1$$

We can then calculate the isotope hit position at a RP and the acceptance

Hit position:

$$x_{RP} = D_{x}(-R_{Rel}) = D_{x}(1 - x_{L})$$

Minimum allowed hit position:

$$x_{min} = 10\sigma_x = 10\sqrt{\beta_x \varepsilon_x + D_x^2 \sigma_p^2}$$

Accelerator Parameters:

$$\varepsilon_{\rm \chi} = 43.2~nm~{\rm (EIC~CDR~Table~3.5)}$$

$$\sigma_p = 6.2 \times 10^{-4}~{\rm (EIC~CDR~Table~3.5)}$$

IR6 Parameters at first RP:

$$\beta_x = 865 m$$

$$D_x = -16.7 cm$$

$$\rightarrow x_{min}^{RP1} = 6.11 cm$$

IR8 Parameters at first RP:

$$\beta_{x} = 2.28 m$$

$$D_{x} = 38.2 cm$$

$$\Rightarrow x_{min}^{RP1} = 0.39 cm$$

We can then calculate the isotope hit position at a RP and the acceptance

Hit position:

$$x_{RP} = D_{x}(-R_{Rel}) = D_{x}(1 - x_{L})$$

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Accelerator Parameters:

$$\varepsilon_{\chi}=43.2~nm$$
 (EIC CDR Table 3.5)
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IR6 Parameters at first RP:

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$$D_x = -16.7 cm$$

$$\rightarrow x_{min}^{RP1} = 6.11 cm$$

IR8 Parameters at first RP:

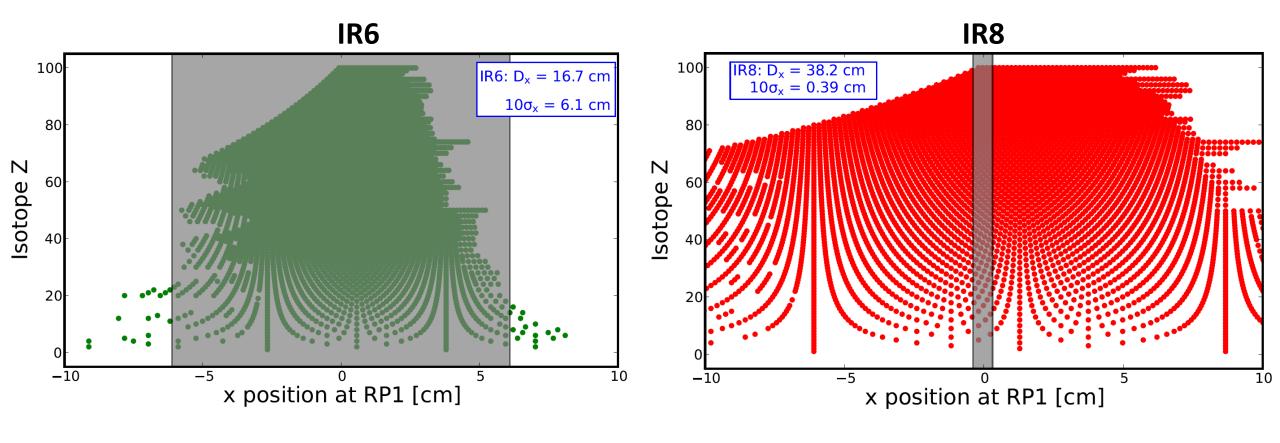
$$\beta_x = 2.28 m$$

$$D_x = 38.2 cm$$

$$\rightarrow x_{min}^{RP1} = 0.39 cm$$

Big acceptance
difference
between the two
IRs is caused by
the second focus
at the RPs in the
IR8 design

Isotope hit positions at the first RP vs. isotope Z

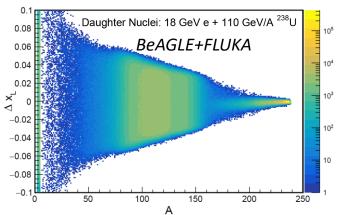


Each point is an individual isotope. All known and potential isotopes which come from a combined NNDC and LISE++ database are included.

Assuming a RP position resolution of 10-100 microns, isotopes with the same Z are well separated.

Some comments on above results

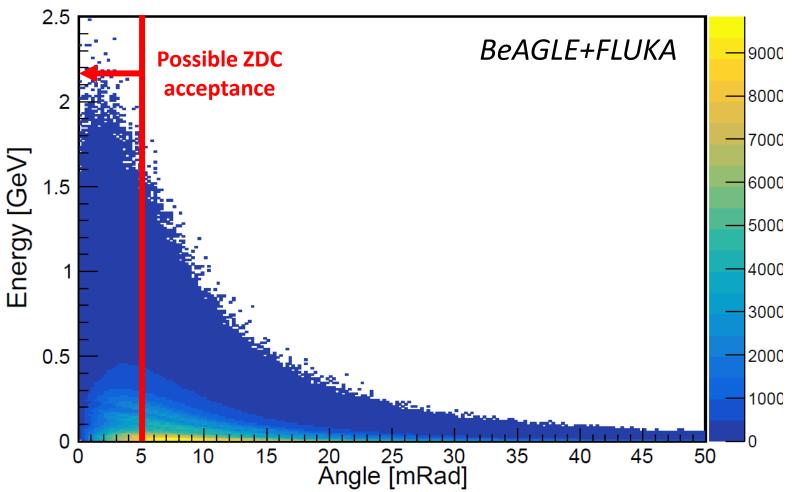
- ☐The above plots show that a large fraction of the potential isotopes can be accepted and correctly identified using the position at the RP.
- ☐ This assumes that the charge of the isotope (Z) has already been determined.
- □ A thin (few mm thick) quartz bar can be placed in the RPs at the second focus behind the tracker to determine Z². The quartz bar would be perpendicular to the beam, extended along the dispersive (x) direction. The number of Cherenkov photons produced will be quite large, and the challenge will be to measure the photons with high enough precision.
- ☐ The assumption made in the plots on the previous slide is that the outgoing isotope has the same momentum-per-nucleon as the ion beam and no angle with respect to the ion beam. Under this assumption, the isotope hit position at the RPs is just a function of (A/Z).
- □ This assumption works well in the evaporation region; but less well in the fission region where the kinetic energy in the excited residual nucleus rest frame can be 1 MeV/nucleon.



$$\Delta x_L = x_L - \left[\frac{\left(\frac{A}{Z}\right)}{\left(\frac{A_{beam}}{Z_{beam}}\right)} \right]$$

Detection of the de-excitation photons

De-excitation Photons: 18 GeV e + 110 GeV/A ²³⁸U



Summary

- ☐ We have shown that the EIC has the potential to produce exotic nuclei.
- ☐ These nuclei can be detected and identified using the proposed optics of the second interaction point with its secondary focus.
- ☐ Studying the level structure of the produced isotopes will be possible through the detection of the de-excitation photons.